

BLOW UP FOR THE 2D EULER EQUATION ON SOME BOUNDED DOMAINS

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ABSTRACT. We find a smooth solution of the 2D Euler equation on a bounded domain which exists and is unique in a natural class locally in time, but blows up in finite time in the sense of its vorticity losing continuity. The domain's boundary is smooth except at two points, which are interior cusps.

1. INTRODUCTION

Consider the 2D Euler equation on a bounded domain $D \subseteq \mathbb{R}^2$, in vorticity formulation:

$$\omega_t + u \cdot \nabla \omega = 0, \quad \omega(0, \cdot) = \omega_0. \quad (1.1)$$

The incompressible velocity u satisfies the no-flow condition $u \cdot n = 0$ on ∂D (with n the unit outer normal to ∂D), and is found from the vorticity ω via the Biot-Savart law

$$u(t, x) = \int_D K_D(x, y) \omega(t, y) dy, \quad (1.2)$$

where the Kernel K_D is given by $K_D(x, y) := \nabla^\perp G_D(x, y)$, with $\nabla^\perp = (-\partial_{x_2}, \partial_{x_1})$ and $G_D \geq 0$ the Green's function for D .

If the domain D and the initial data ω_0 are smooth, results on existence, uniqueness, and global regularity of solutions of the 2D Euler equation go back to the 1930s works of Wolibner [13] and Hölder [4]. Even though the solutions in this case are globally regular, infinite-time growth of their derivatives and small scale creation appear to be ubiquitous in two dimensional fluids.

The best known upper bound on the growth of the gradient of the vorticity is double-exponential in time (see [16], even though the result is also implicit in earlier works). The question of how fast it can actually grow has received much attention in the literature. First results of this type were due to Yudovich [14, 15], who constructed examples with unbounded growth of the vorticity gradient on the boundary of the domain. Recently, such growth at the boundary was shown to be fairly generic in [11] and [6]. Better quantitative estimates were obtained by Nadirashvili [12], who built examples with at least linear growth, and Denisov [1], whose construction gives a super-linear lower bound for the vorticity gradient. Denisov also constructed examples with a double-exponential growth for an arbitrary but finite time [2]. In these examples, growth happens in the bulk of the fluid. Very recently, Kiselev and Šverák [5] constructed examples with a double-exponential in time growth of the vorticity gradient on the boundary, which is of course the fastest possible rate of growth. Whether such growth can happen in the bulk remains open. The best result here so far is due to Zlatoš [18], who built examples with at least exponential in time growth.

The 2D Euler equation in less regular domains and with less regular initial data has also been studied extensively. Yudovich theory [17] (see also [9, 10]) guarantees existence and uniqueness of a bounded weak solution ω if the domain is C^∞ and the initial data ω_0 is L^∞ . Uniqueness comes from the fact that the fluid velocity corresponding to a bounded vorticity ω is log-Lipshitz. This makes bounded vorticities a very natural class to consider (although incremental generalizations are possible). The smoothness assumptions on the boundary of the domain can also be weakened, and Yudovich theory can be extended to domains which are only $C^{1,1}$ [3]. Moreover, if the initial data and domain have a higher regularity, the solution inherits it as well [10]. The 2D Euler equation has also been studied on certain types of domains with singular points, including domains with corners. We refer to recent works [3, 7, 8] for some results on existence, uniqueness, and regularity of solutions, as well as for further references.

Our main purpose in this note is to provide an example showing that if the boundary of the domain is not *everywhere* sufficiently regular, then much more dramatic effects than just a fast growth of the derivatives of solutions can take place. We show that there are solutions of the 2D Euler equation on domains with two cusps (and smooth elsewhere) which are smooth locally in time but blow up in finite time, in the sense of the vorticity losing continuity. We note that our construction will rely only on the above existence and uniqueness results for smooth domains and initial data, due to a special structure of our example.

The example involves a “double stadium” domain. Let

$$\tilde{D}_r^\pm := [(-r, r) \times (\pm 1 - r, \pm 1 + r)] \cup B_r(-r, \pm 1) \cup B_r(r, \pm 1)$$

be the “stadium” domain in \mathbb{R}^2 with width $r > 0$ and centered at $(0, \pm 1)$. While \tilde{D}_r^\pm are defined explicitly, they are only $C^{1,1}$. Let us therefore define domains D_r^\pm by making the connections between the circular arcs and the horizontal intervals (which form $\partial\tilde{D}_r^\pm$) infinitely smooth. This can be done by adjusting the circular arcs near the points of contact with the intervals, and such an adjustment can be made as small as we want in C^1 . Let us do this so that again all the D_r^\pm are rescaled copies of each other, centered at $(0, \pm 1)$. Finally, let

$$D' := D_1^+ \cup D_1^- \cup [(-1, 1) \times \{0\}]$$

and denote $D := D_1^+$ the upper half of D' .

If $\omega_0 \in L^\infty(D')$ is odd in x_2 , then the Yudovich theory shows that there is a unique odd-in- x_2 weak solution ω to (1.1) on $D' \times \mathbb{R}$, namely the one whose restriction to $D \times \mathbb{R}$ solves (1.1) there. Moreover, if $\omega_0 \in C(D')$, then it follows from Yudovich theory that $\omega(t, \cdot) \in C(D)$ for all $t > 0$. The following theorem provides a necessary and sufficient condition for ω to remain continuous on the whole domain D' .

Theorem 1.1. *Let $\omega_0 \in C(\bar{D}')$ be non-negative on D and odd in x_2 . Then the unique odd-in- x_2 solution ω to (1.1) on $D' \times \mathbb{R}$ is continuous for all $t > 0$ if and only if $\omega_0(\partial D) = 0$.*

Remark. If in addition we have $\omega_0 \in C^\infty(\bar{D})$, then classical regularity theory gives us $\omega(t, \cdot) \in C^\infty(D)$ for all $t > 0$. If furthermore ω_0 is supported away from $[-1, 1] \times \{0\}$, then finite speed of propagation (i.e., bounded u) implies $\omega(t, \cdot) \in C^\infty(D')$ for all small enough $t > 0$. By Theorem 1.1, such a solution will become discontinuous in finite time unless ω_0 vanishes on all of ∂D .

2. PROOF OF THEOREM 1.1

If $\omega_0(\partial D) = 0$, then the result follows from Yudovich theory because both D and ∂D are invariant under the flow map generated by u .

Let us now assume that $\omega_0(x_0) > 0$ for some $x_0 \in \partial D$. We will show that if $X'(t; x) = u(t, X(t; x))$ and $X(0; x) = x$ for $x \in \bar{D}$, then there is $T > 0$ such that $X(T; x_0) = 0$. This, the continuity of ω_0 and of $X(t; \cdot)$ on \bar{D} , and oddness of ω in x_2 then show that $\omega(T, \cdot)$ is discontinuous at the origin.

Let $\delta > 0$ be such that $|\omega_0^{-1}([\delta, \infty))| \geq 30\delta (> (4 + 2\pi)2\delta + \delta)$. Because u is incompressible and $\omega(t, \cdot) \geq 0$ on D , it follows that for each $t \geq 0$,

$$|\{y \in D_{1-2\delta}^+ \mid \omega(t, y) \geq \delta\}| \geq \delta. \quad (2.1)$$

Next let

$$\kappa := \inf_{x \in \partial D_{1-\delta}^+ \text{ \& } y \in \bar{D}_{1-2\delta}^+} G_D(x, y).$$

Continuity of G_D away from the diagonal $x = y$, positivity of G_D on D , and compactness of \bar{D}_r^+ show that $\kappa > 0$. Let $v > 0$ solve $\Delta v = 0$ on $D \setminus \bar{D}_{1-\delta}^+$ with $v = 0$ on ∂D and $v = \kappa$ on $\partial D_{1-\delta}^+$. By the Hopf Lemma it follows that $\varepsilon := \inf_{x \in \partial D} |\nabla v(x)| > 0$.

We then have $G_D(x, y) \geq v(x)$ for $x \in D \setminus \bar{D}_{1-\delta}^+$ and $y \in \bar{D}_{1-2\delta}^+$, because the inequality holds for $x \in \partial(D \setminus \bar{D}_{1-\delta}^+)$ (and both functions are harmonic in x). So we have $|K_D(x, y)| \geq \varepsilon$ for $x \in \partial D$ and $y \in \bar{D}_{1-2\delta}^+$. This and (2.1) show that for $x \in \partial D$, the right-hand side of (1.2) equals $a(t, x)(n_2(x), -n_1(x))$, with $a(t, x) \geq \varepsilon\delta^2$. Thus $X(t; x_0)$ travels clockwise along ∂D with velocity $\geq \varepsilon\delta^2$, which means that $X(T; x_0) = 0$ for some $T > 0$. The proof is finished.

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